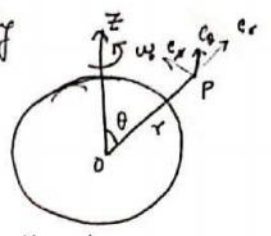


Steady Motion of viscous Fluid Due to a Slowly Rotating sphere

For the present problem, the velocity components are

Let sphere is rotating slowly with angular velocity ω_0 about the fixed axis $\theta=0$ passing through the centre of sphere. Due to viscosity surrounding fluid will also rotate in the same direction. Thus there will be only motion in the direction



of \hat{e}_θ and $V_r = V_\theta = 0$. Let ω be the angular velocity at any point P. This is to be noted that angular velocity ω is a fn of r only. In Cartesian Co-ordinates the components of velocity at P are

$$u = -\omega y, \quad v = \omega x, \quad w = 0$$

$$\vec{v} = \omega \hat{k} \times (\vec{r}) = \omega \hat{k} \times (x\hat{i} + y\hat{j} + z\hat{k})$$

$$\vec{v} = \vec{\omega} \times \vec{r}$$

$$= \omega \hat{k} \times (x\hat{i} + y\hat{j} + z\hat{k})$$

$$u\hat{i} + v\hat{j} + w\hat{k} = \omega x\hat{j} - \omega y\hat{i}$$

i.e. $u = -\omega y, \quad v = \omega x, \quad w = 0$ — (1)

The Navier-Stokes equations in the absence of the body forces is

$$\rho \frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \nabla) \vec{v} = -\nabla p + \mu \nabla^2 \vec{v} \quad \text{--- (2)}$$

Due to steady motion $\frac{\partial \vec{v}}{\partial t} = 0$. Again, since \vec{v} is very small, so neglecting squares of velocities, we have $(\vec{v} \cdot \nabla) \vec{v} = 0$, Hence (2) reduces to

$$\mu \nabla^2 \vec{v} = \nabla p$$

~~$$\mu [\hat{i} \nabla^2 u + \hat{j} \nabla^2 v + \hat{k} \nabla^2 w] = \frac{\partial p}{\partial x} \hat{i} + \frac{\partial p}{\partial y} \hat{j} + \frac{\partial p}{\partial z} \hat{k}$$~~

$$\mu [\hat{i} \nabla^2 u + \hat{j} \nabla^2 v] = \frac{\partial p}{\partial x} \hat{i} + \frac{\partial p}{\partial y} \hat{j} + \frac{\partial p}{\partial z} \hat{k}$$

i.e.
$$\left. \begin{aligned} \mu \nabla^2 u &= \frac{\partial p}{\partial x} \\ \mu \nabla^2 v &= \frac{\partial p}{\partial y} \\ 0 &= \frac{\partial p}{\partial z} \end{aligned} \right\} \text{--- (3)}$$

Now we find $\nabla^2 u$ in term of the derivatives w.r.t r . For this we will evaluate

$$\frac{\partial u}{\partial x} = \frac{\partial}{\partial x}(-y\omega) = -y \frac{\partial \omega}{\partial x}$$

$$\frac{\partial^2 u}{\partial x^2} = -y \frac{\partial^2 \omega}{\partial x^2}$$

and $\frac{\partial u}{\partial y} = \frac{\partial}{\partial y}(-y\omega) = -\omega - y \frac{\partial \omega}{\partial y}$

$$\frac{\partial^2 u}{\partial y^2} = -\frac{\partial \omega}{\partial y} - \left(y \frac{\partial^2 \omega}{\partial y^2} + \frac{\partial \omega}{\partial y} \right) = -y \frac{\partial^2 \omega}{\partial y^2} - 2 \frac{\partial \omega}{\partial y}$$

and $\frac{\partial u}{\partial z} = -y \frac{\partial \omega}{\partial z}$, $\frac{\partial^2 u}{\partial z^2} = -y \frac{\partial^2 \omega}{\partial z^2}$

thus $\nabla^2 u = \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2}$

$$= -y \left(\frac{\partial^2 \omega}{\partial x^2} + \frac{\partial^2 \omega}{\partial y^2} + \frac{\partial^2 \omega}{\partial z^2} \right) - 2 \frac{\partial \omega}{\partial y}$$

$$= -y \left[\frac{\partial^2 \omega}{\partial x^2} + \frac{\partial^2 \omega}{\partial y^2} + \frac{\partial^2 \omega}{\partial z^2} + \frac{2}{y} \frac{\partial \omega}{\partial y} \right] \quad \text{--- (4)}$$

Now since $\frac{\partial \omega}{\partial x} = \frac{\partial \omega}{\partial r} \cdot \frac{r}{y} = \frac{r}{y} \frac{\partial \omega}{\partial r}$

$$\frac{1}{y} \frac{\partial \omega}{\partial y} = \frac{1}{r} \frac{d\omega}{dr}$$

Now since $\frac{\partial \omega}{\partial x} = \frac{\partial \omega}{\partial r} \cdot \frac{r}{x} = \frac{r}{x} \frac{\partial \omega}{\partial r} = \frac{r}{y} \frac{d\omega}{dr}$

$$\frac{1}{x} \frac{\partial \omega}{\partial x} = \frac{1}{r} \frac{d\omega}{dr} \quad \text{--- (5)}$$

and $\frac{\partial^2 \omega}{\partial x^2} = \frac{\partial}{\partial x} \left(\frac{\partial \omega}{\partial x} \right) = \frac{\partial}{\partial x} \left(\frac{r}{y} \frac{d\omega}{dr} \right)$

$$= \frac{1}{y} \frac{d\omega}{dr} + x \frac{\partial}{\partial x} \left(\frac{1}{r} \frac{d\omega}{dr} \right)$$

$$= \frac{1}{y} \frac{d\omega}{dr} + x \cdot \frac{\partial}{\partial r} \left(\frac{1}{r} \frac{d\omega}{dr} \right) \frac{\partial x}{\partial r}$$

$$= \frac{1}{y} \frac{d\omega}{dr} + \frac{x^2}{y} \cdot \left(\frac{1}{r} \frac{d^2 \omega}{dr^2} - \frac{1}{r^2} \frac{d\omega}{dr} \right)$$

$$= \frac{x^2}{r^2} \frac{d^2 \omega}{dr^2} + \frac{x^2 - r^2}{r^3} \frac{d\omega}{dr} \quad \text{--- (6)}$$



Similarly

$$\frac{\partial^2 w}{\partial y^2} = \frac{y^2}{r^2} \frac{d^2 w}{dr^2} + \frac{r^2 y^2}{r^2} \frac{dw}{dr} \quad \text{--- (7)}$$

$$\text{and } \frac{\partial^2 w}{\partial z^2} = \frac{z^2}{r^2} \frac{d^2 w}{dr^2} + \frac{r^2 z^2}{r^2} \frac{dw}{dr} \quad \text{--- (8)}$$

Adding (6), (7) and (8), we have

$$\begin{aligned} \frac{\partial^2 w}{\partial x^2} + \frac{\partial^2 w}{\partial y^2} + \frac{\partial^2 w}{\partial z^2} &= \frac{dw}{dr} + \frac{2r^2 z^2}{r^2} \frac{dw}{dr} \\ &= \frac{d^2 w}{dr^2} + \frac{2}{r} \frac{dw}{dr} \quad \text{--- (9)} \end{aligned}$$

Using (5) and (9), (4) reduces to

~~$$\begin{aligned} \text{Now, } \nabla^2 u &= \nabla^2 (yw) = - \nabla^2 (yw) \\ &= - \left[\frac{\partial^2}{\partial x^2} (yw) + \frac{\partial^2}{\partial y^2} (yw) + \frac{\partial^2}{\partial z^2} (yw) \right] \\ &= - \left[y \frac{\partial^2 w}{\partial x^2} + 2 \frac{\partial}{\partial y} (w + y \frac{\partial w}{\partial y}) + y \frac{\partial^2 w}{\partial z^2} \right] \\ &= - \left[y \frac{\partial^2 w}{\partial x^2} + \left(\frac{\partial w}{\partial y} + y \frac{\partial^2 w}{\partial y^2} + \frac{\partial w}{\partial y} \right) + y \frac{\partial^2 w}{\partial z^2} \right] \end{aligned}$$~~

~~$$\text{By (4) } \nabla^2 u = - \left[y \left(\frac{\partial^2 w}{\partial x^2} + \frac{\partial^2 w}{\partial y^2} + \frac{\partial^2 w}{\partial z^2} \right) + 2 \frac{\partial w}{\partial y} \right]$$~~

~~$$\nabla^2 u = - \left[y \cdot \left(\frac{dw}{dr} + \frac{2}{r} \frac{dw}{dr} \right) + 2 \frac{y}{r} \frac{dw}{dr} \right]$$~~

~~$$\nabla^2 u = - y \left[\frac{d^2 w}{dr^2} + \frac{4}{r} \frac{dw}{dr} \right] \quad \text{--- (10)}$$~~

~~$$\text{Similarly } \nabla^2 v = x \left[\frac{d^2 w}{dr^2} + \frac{4}{r} \frac{dw}{dr} \right] \quad \text{--- (11)}$$~~

with these expressions for $\nabla^2 u$ and $\nabla^2 v$ Eqⁿ (3) reduces to be

$$\left. \begin{aligned} - \mu y \left(\frac{dw}{dr} + \frac{4}{r} \frac{dw}{dr} \right) &= \frac{\partial p}{\partial x} \\ \text{and } \mu x \left(\frac{dw}{dr} + \frac{4}{r} \frac{dw}{dr} \right) &= \frac{\partial p}{\partial y} \\ 0 &= \frac{\partial p}{\partial z} \end{aligned} \right\} \text{--- (12)}$$



All equations in Eq. (12) are satisfied by taking $p = \text{constant}$ and

$$\frac{d^2\omega}{dr^2} + \frac{4}{r} \frac{d\omega}{dr} = 0$$

$$\Rightarrow r^4 \frac{d^2\omega}{dr^2} + 4r^3 \frac{d\omega}{dr} = 0$$

$$\Rightarrow \frac{d}{dr} \left(r^4 \frac{d\omega}{dr} \right) = 0 \quad \text{--- (13)}$$

Integrating (13),

$$r^4 \frac{d\omega}{dr} = C_1$$

$$\frac{d\omega}{dr} = \frac{C_1}{r^4}$$

$$\omega = -\frac{C_1}{3r^3} + C_2$$

$$\omega = \frac{C_3}{r^3} + C_2 \quad \text{--- (14)}$$

If motion is generated by a solid sphere of radius 'a' rotating with angular velocity ω_0 . we have B.C.

$$(i) \omega = 0 \quad \text{at } r = 0$$

$$(ii) \omega = \omega_0 \quad \text{at } r = a.$$

using these B.C. in (14) we get

$$C_2 = 0$$

$$\omega C_3 = \omega_0 \cdot a^3$$

$$\text{Thus } \omega = \frac{a^3 \omega_0}{r^3} \quad \text{--- (15)}$$

* when outer sphere is stationary and is of radius b. then B.C. (i) will be $\omega = 0$ at $r = b$. we have

$$C_2 = -\frac{\omega_0 a^3}{(b^3 - a^3)}$$

$$C_3 = \frac{\omega_0 a^3 b^3}{(b^3 - a^3)}$$

and
$$\omega = \frac{\omega_0 a^3}{(b^3 - a^3)} \left(\frac{b^3}{r^3} - 1 \right)$$

If motion is due to rotation of two concentric spheres of radius a and b with the angular velocity ω_1



~~Here we have~~
 velocity of fluid Particle in spherical Polar Coordinate is
 $\vec{V} = \dot{r} \hat{e}_r + r \dot{\theta} \hat{e}_\theta + r \sin\theta \dot{\phi} \hat{e}_\phi$
 here $V_r = V_\theta = 0$ and $V_\phi = \omega r \sin\theta$

The only stress component is

$$P_{r\phi} = \mu \left[\frac{1}{r \sin\theta} \frac{\partial V_r}{\partial \phi} + r \frac{\partial}{\partial r} \left(\frac{V_\phi}{r} \right) \right]$$

$$V_r = 0, V_\phi = \omega r \sin\theta$$

$$P_{r\phi} = \mu r \cdot \frac{\partial}{\partial r} (\omega \sin\theta) = \mu \sin\theta \cdot \frac{\partial}{\partial r} \left(\frac{a^3}{r^2} \omega \right)$$

$$= \mu \sin\theta \cdot a^3 \omega \left(-\frac{3}{r^3} \right)$$

$$P_{r\phi} = -\frac{3\mu a^3 \omega \sin\theta}{r^3}$$

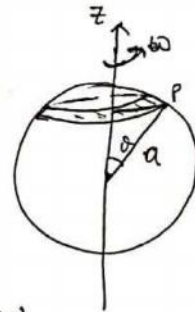
Stress on the boundary of sphere

$$(P_{r\phi})_{r=a} = -3\mu \omega \sin\theta$$

Couple on the sphere

$$= \int_0^{2\pi} (\text{Moment about the } z \text{ axis}) d\theta$$

$$(2\pi a \sin\theta) (a d\theta)$$



$$= \int_0^{2\pi} (P_{r\phi} \cdot a \sin\theta) \cdot 2\pi a^2 \sin\theta d\theta$$

$$= \int_0^{2\pi} -3\mu \omega \sin\theta \cdot a \sin\theta \cdot 2\pi a^2 \sin\theta d\theta$$

$$= -6\mu\pi \omega a^3 \int_0^{2\pi} \sin^3 \theta d\theta = -6\mu\pi \omega a^3 \cdot 2 \int_0^{\pi/2} \sin^3 \theta d\theta$$



$$-12\mu\pi \omega a^3 \cdot \frac{2}{3} = -8\mu\pi \omega a^3$$